

# *The static three-quark potential from the Polyakov loop correlation function*

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Yoshiaki Koma (Numazu College of Technology)  
Miho Koma (Nihon University)

— Lattice 2014 @ Columbia U —



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# Introduction

## ► Heavy quarkonium: $Q\bar{Q}$

⇒ potential NRQCD [cf. Brambilla et al.(RMP77)]

$$H = \frac{\vec{p}_1^2}{2m_1} + \frac{\vec{p}_2^2}{2m_2} + \underbrace{V^{(0)}(x_1, x_2)}_{\text{well known}} + \underbrace{O(\frac{1}{m})}_{\text{known}} \quad [\text{Koma et al.(PRL97)ff}]$$

⇒ quarkonium can be studied “quantum mechanically”

## ► Heavy baryon: $QQQ$

$$H = \frac{\vec{p}_1^2}{2m_1} + \frac{\vec{p}_2^2}{2m_2} + \frac{\vec{p}_3^2}{2m_3} + \underbrace{V^{(0)}(x_1, x_2, x_3)}_{\text{known*}} + \underbrace{O(\frac{1}{m})}_{\text{unknown}}$$

⇒ precise functional form of the three-quark potential ?

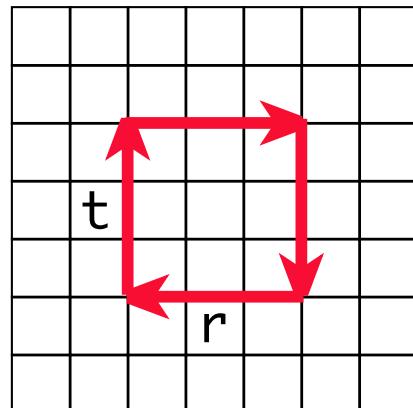
\*[cf. Tahakashi et al.(PRD65), Alexandrou et al.(PRD65)]

# How to compute the potential

- ▶ Inter-quark potential can be extracted from the Wilson loop or Polyakov loop correlation function (PLCF)

## Wilson loop

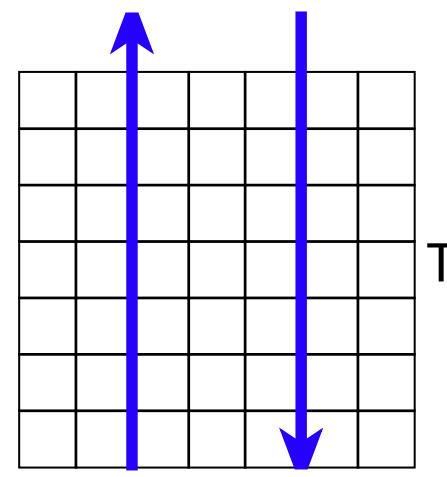
$$V^{(0)}(r) = - \lim_{t \rightarrow \infty} \frac{1}{t} \ln \langle W(r, t) \rangle$$



$W(r, t)$

## PLCF

$$V^{(0)}(r) = - \lim_{T \rightarrow \infty} \frac{1}{T} \ln \langle P(0)P(r)^* \rangle$$



$P(0)P(r)^*$

(popular at finite temperature)

# Wilson loop vs. PLCF

►  $W(r, t) = \mathbb{L}(r, 0)_{\alpha\gamma}^* \{\mathbb{T}(0)\mathbb{T}(a)\cdots\mathbb{T}(t-a)\}_{\alpha\beta\gamma\delta} \mathbb{L}(r, t)_{\beta\delta}$

$$\Rightarrow \langle W(r, t) \rangle = \sum_{n=0}^{\infty} w_n(r, t) e^{-E_n(r)t} \quad (w_n(r, t) = \langle 0 | \mathbb{L}(r, 0)^* | n(r) \rangle \langle n(r) | \mathbb{L}(r, t) | 0 \rangle)$$

$$\Rightarrow -\frac{1}{t} \ln \langle W(r, t) \rangle = \underbrace{E_0(r)}_{=V^{(0)}} - \underbrace{\frac{1}{t} \ln w_0(r, t) + O(\frac{1}{t} e^{-(E_1-E_0)t})}_{\text{unwanted contributions}}$$

►  $P(0)P(r)^* = \{\mathbb{T}(0)\mathbb{T}(a)\cdots\mathbb{T}(T-a)\}_{\alpha\alpha\gamma\gamma}$

$$\Rightarrow \langle P(0)P(r)^* \rangle = \sum_{n=0}^{\infty} w_n e^{-E_n(r)T} \quad (w_0 = 1)$$

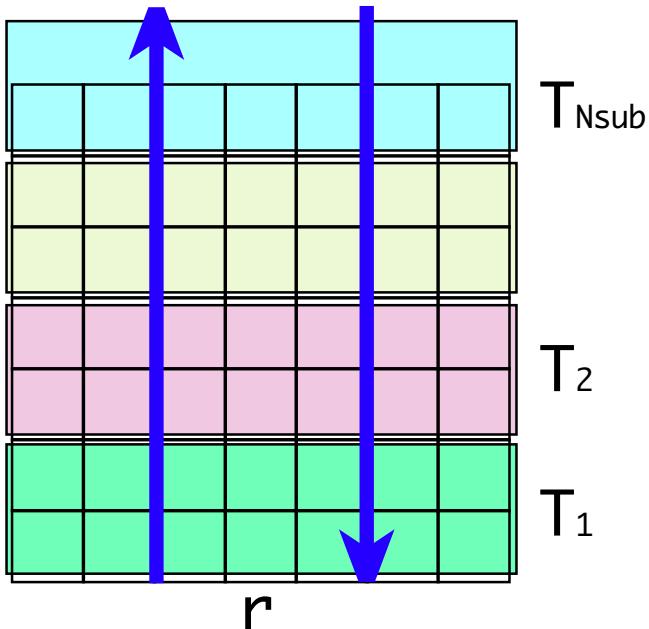
$$\Rightarrow -\frac{1}{T} \ln \langle P(0)P(r)^* \rangle = \underbrace{E_0(r)}_{=V^{(0)}} + \underbrace{O(\frac{1}{T} e^{-(E_1-E_0)T})}_{\text{can be neglected}}$$

# PLCF at zero temperature

- Difficult to obtain accurate PLCFs for larger  $r$  and  $T$ , since the expectation values are exponentially smaller

⇒ use the multilevel algorithm

[Lüscher&Weisz, JHEP0109(2001)010, JHEP0207(2002)049]



$$P(0)P(r)^* = [T_1][T_2] \cdots [T_{N_{sub}}]$$

construct the PLCF form the average of the sublattice correlator  $T(r)$

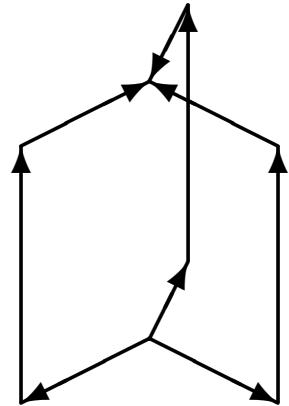
⇒ exponential noize reduction

(e.g.)  $N_{sub} = 4$

$$O(10^{-20}) \leftarrow O(10^{-5}) \cdot O(10^{-5}) \cdot O(10^{-5}) \cdot O(10^{-5})$$

# Three-quark potential

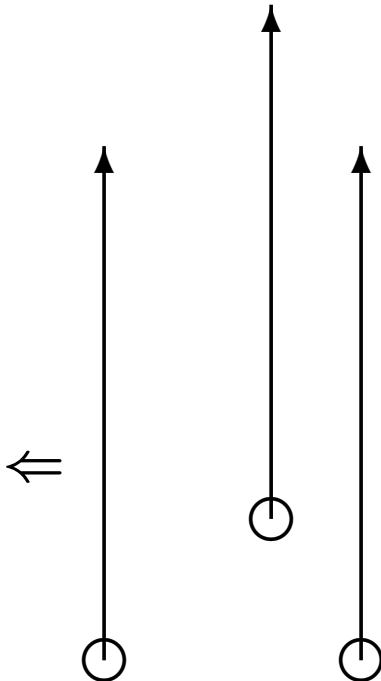
- **Wilson loop** [Takahashi et al.,(PRD65), Alexandrou et al.(PRD65)]



⇒ **junction in source**  
(may cause systematic effects)

- **PLCF**

**NO junction**



$$P(x_1)P(x_2)P(x_3)$$

# Results

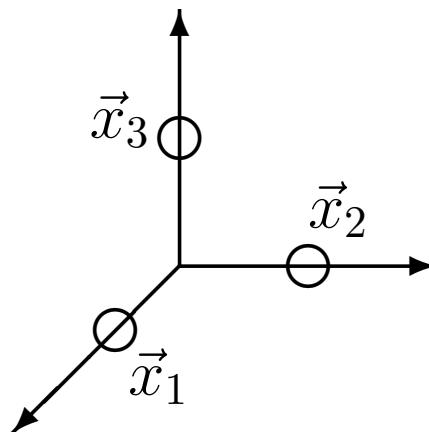
## ► Simulation details

- SU(3) Wilson gauge action

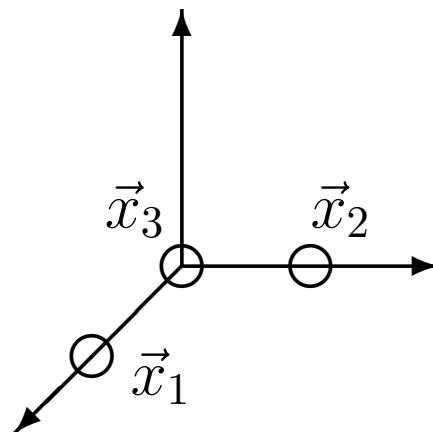
$\beta$	$L^3 T$	$a$ [fm]	$N_{\text{sub}}$	$N_{\text{iupd}}$	$N_{\text{conf}}$
5.85	$24^4$	0.123	8	500000	8
6.00	$24^4$	0.093	6	500000	$1 \sim 8$
6.30	$24^4$	0.059	4	400000	4

[NEC-SX8@RCNP, Osaka Univ]

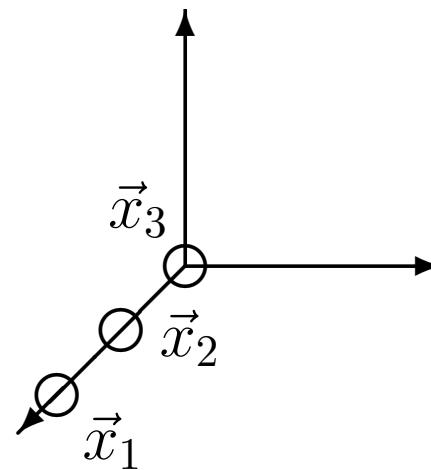
- quark locations



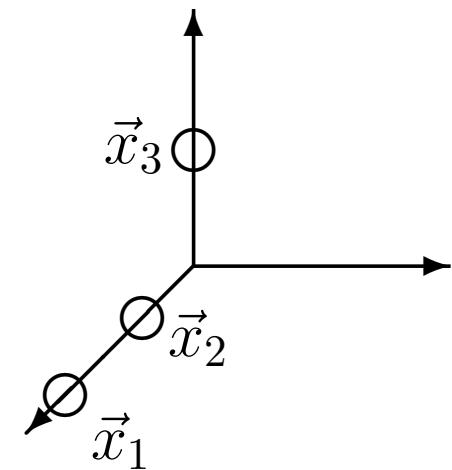
(regular, isosceles)



(right)



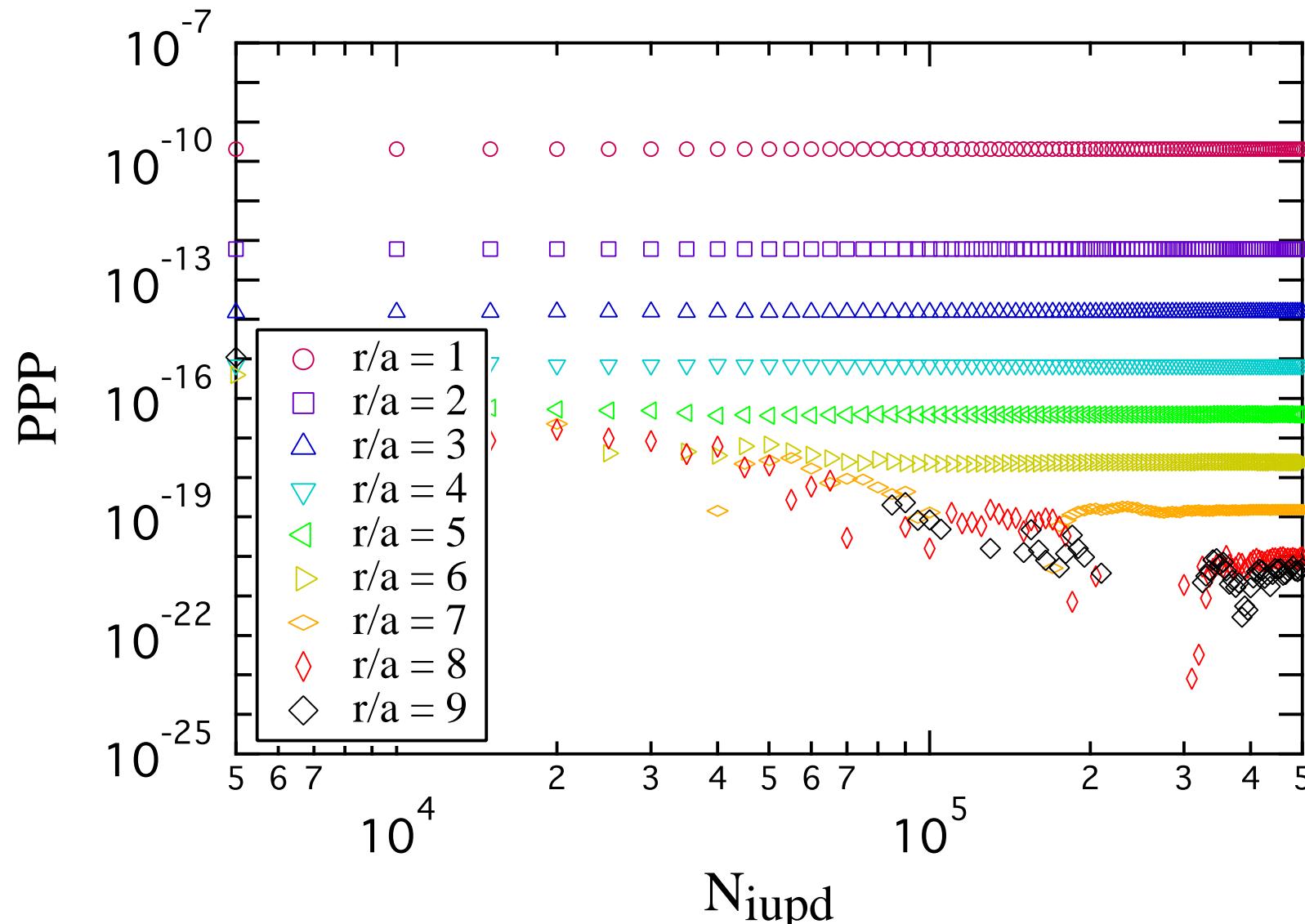
(line)



(obtuse)

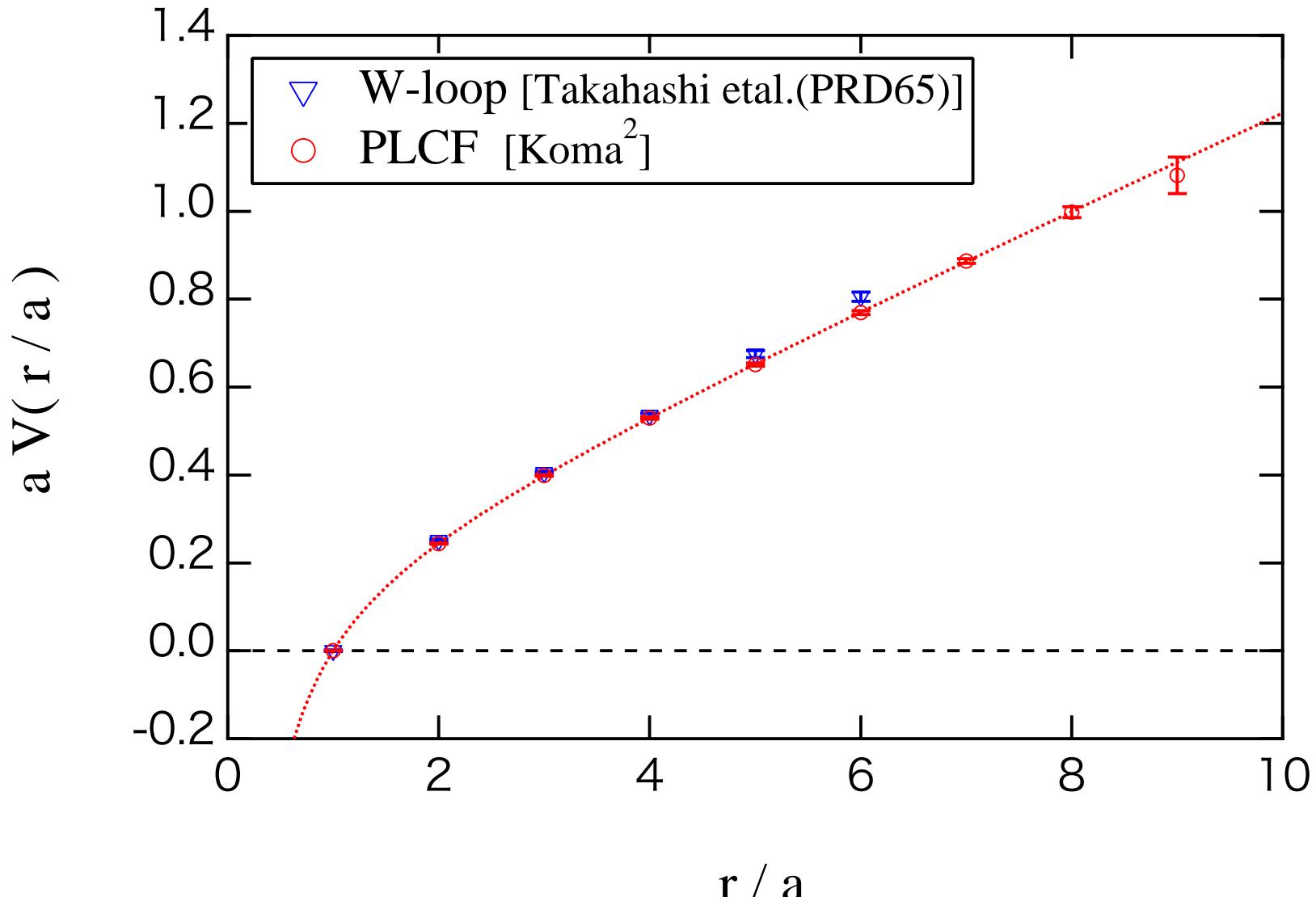
# Three-quark PLCF (regular)

► IUPD history of PLCF for  $\vec{x}_i = r\vec{e}_i$  ( $\beta=6.0$ ,  $24^4$  lattice,  $N_{\text{sub}}=6$ )



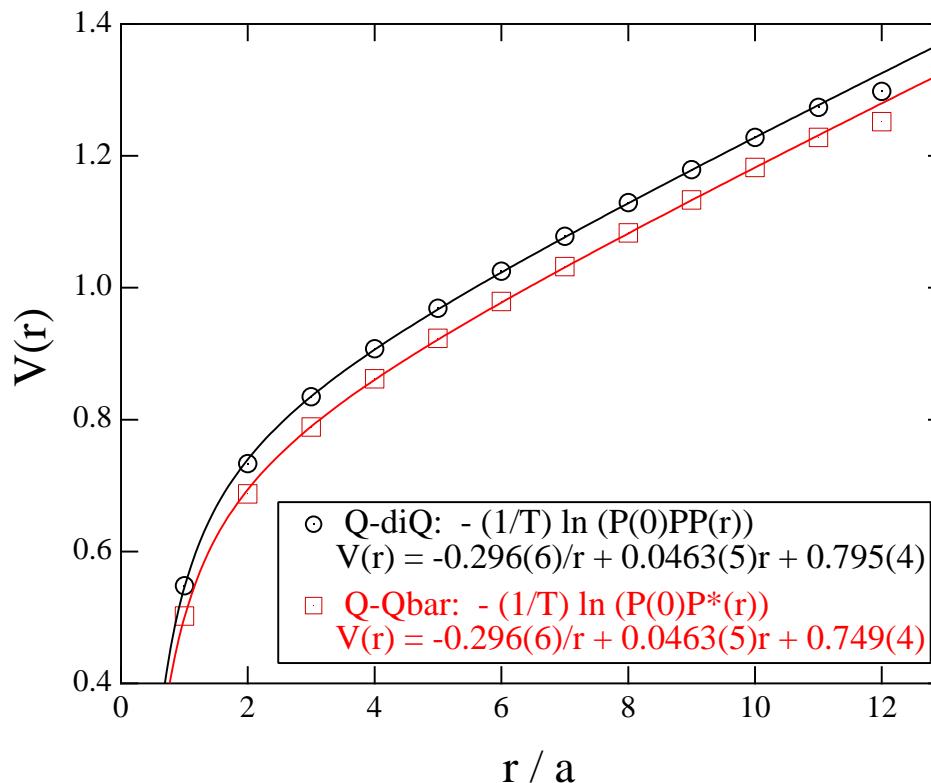
# Three-quark potential (regular)

►  $\vec{x}_i = r \vec{e}_i$  ( $\beta = 6.0$ ,  $24^4$  lattice,  $N_{\text{sub}} = 6$ ,  $N_{\text{iupd}} = 500000$ ,  $N_{\text{conf}} = 8$ )



# Three-quark potential ( $Q-Q\bar{Q}$ system)

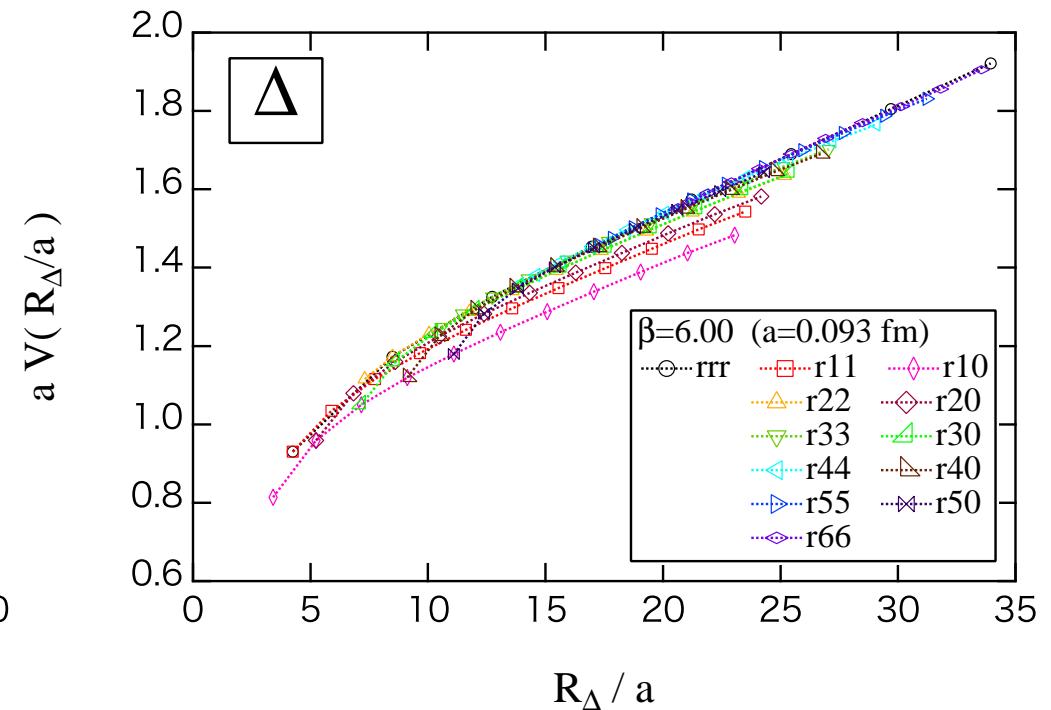
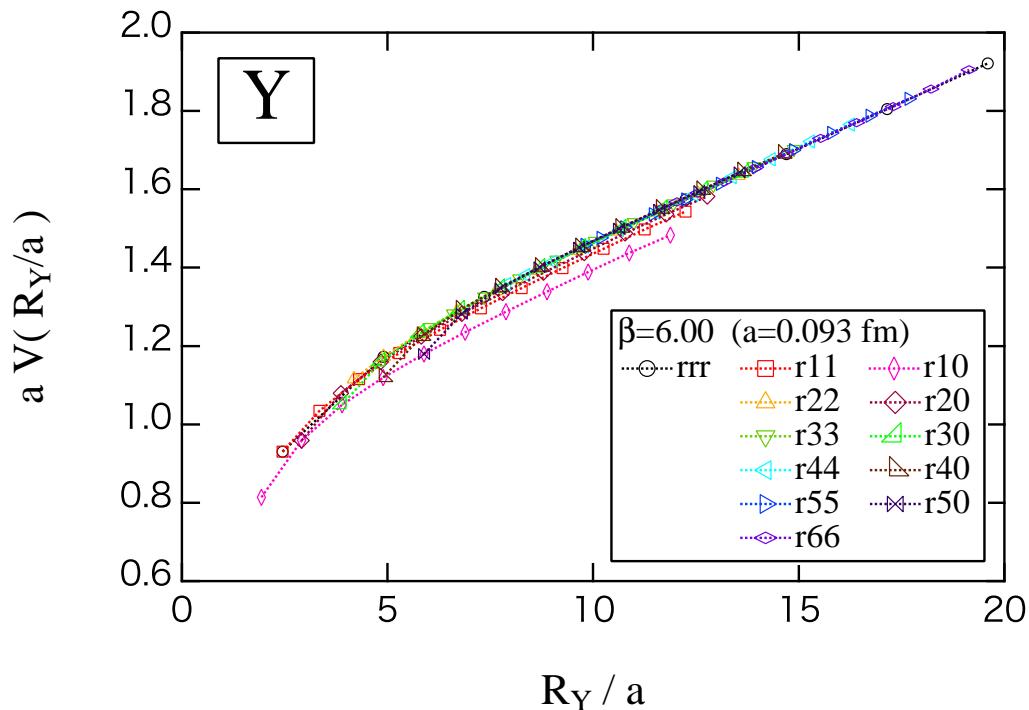
- $\vec{x}_1 = r\vec{e}_1, \vec{x}_2 = \vec{x}_3 = \vec{0}$  ( $\beta = 6.0, 24^4$  lattice,  $N_{\text{sub}} = 6$ )



- $Q-\bar{Q}$  and  $Q-Q\bar{Q}$  potential are the same except constant shift  
 $3 \otimes (\bar{3} \oplus 6) \Rightarrow 3 \otimes \bar{3}$  confirmed [cf. Bissey, Signal & Leinweber, PRD80]

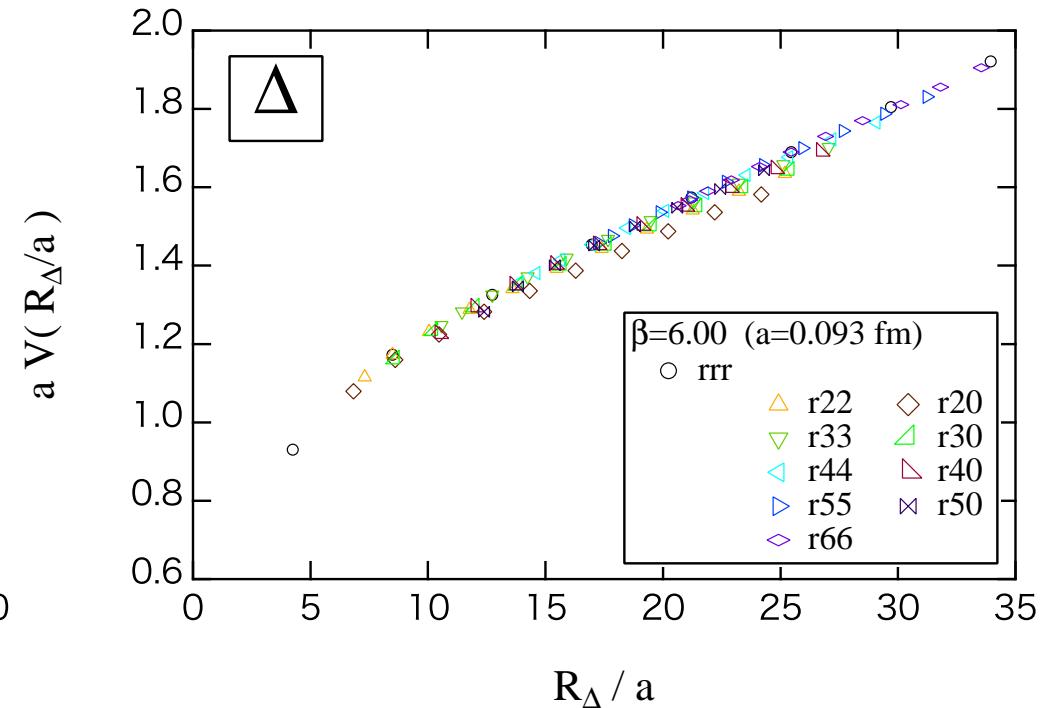
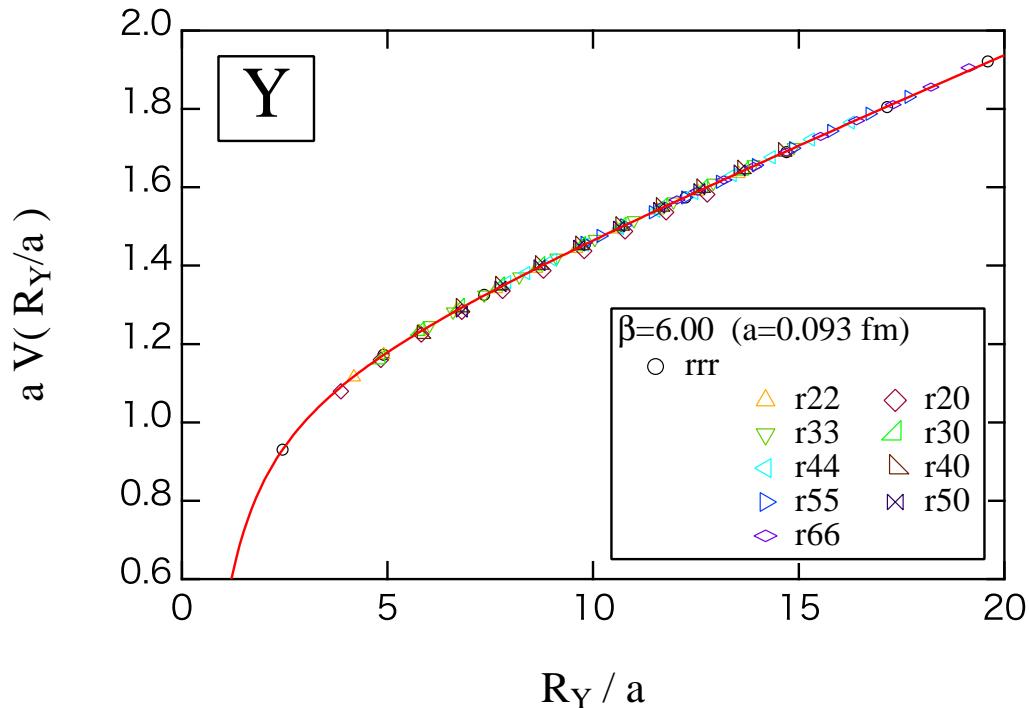
# Three-quark potential (regular, isosceles, right)

- ▶ Various 3-quark locations with  $Y$  and  $\Delta$  parametrization  
( $\beta=6.0$ ,  $24^4$  lattice,  $N_{\text{sub}}=6$ )



# Three-quark potential (regular, isosceles, right)

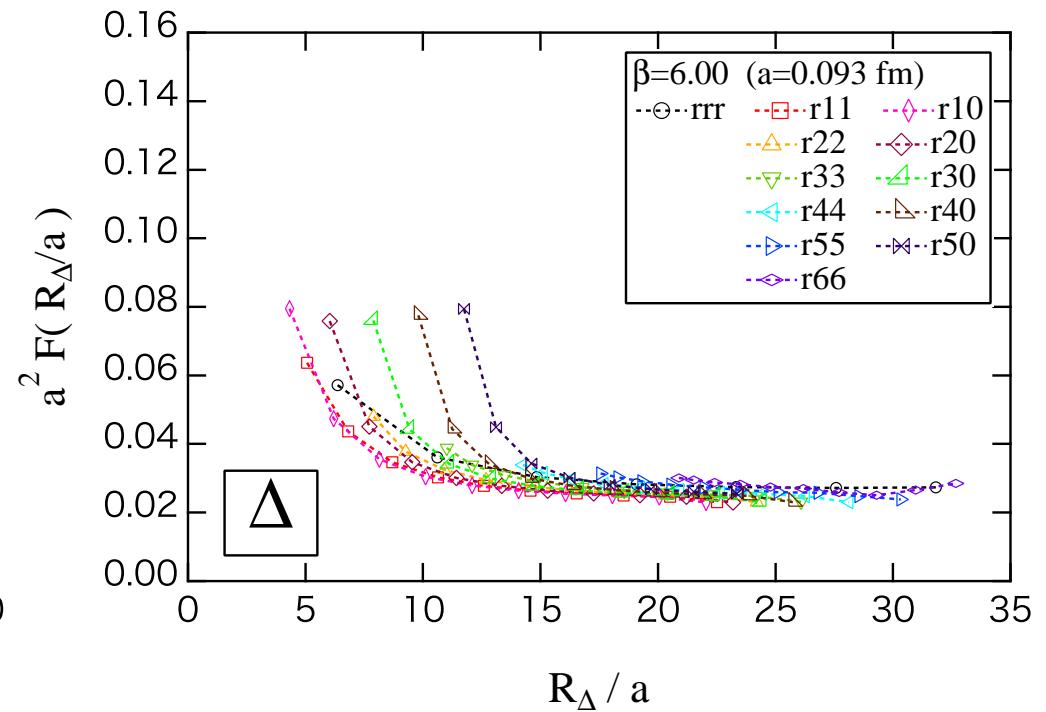
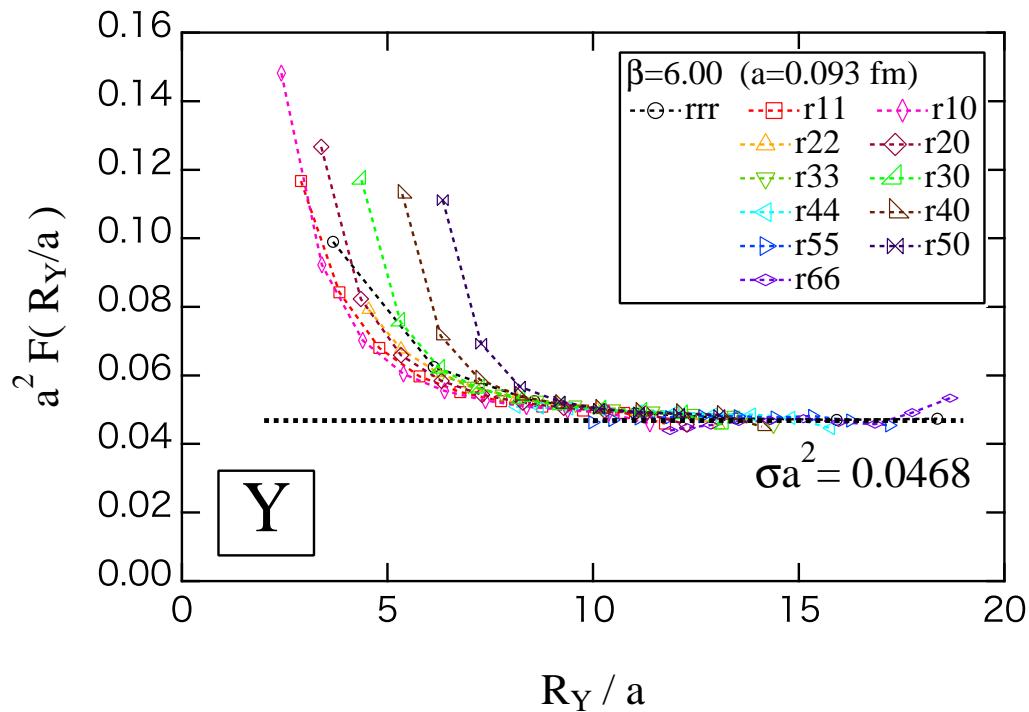
- ▶ Various 3-quark locations with  $Y$  and  $\Delta$  parametrization  
( $\beta=6.0$ ,  $24^4$  lattice,  $N_{\text{sub}}=6$ )



- ▶ The potential with  $Y$  parametrization falls into **ONE CURVE** whenever three quarks are separated with each other

# Three-quark force (regular, isosceles, right)

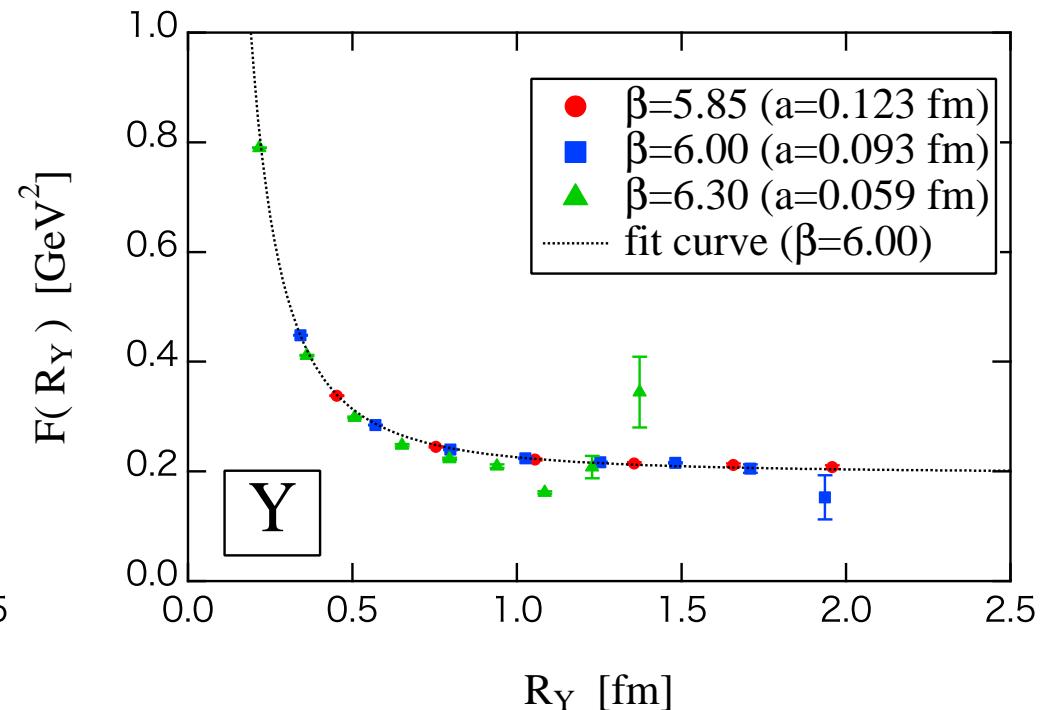
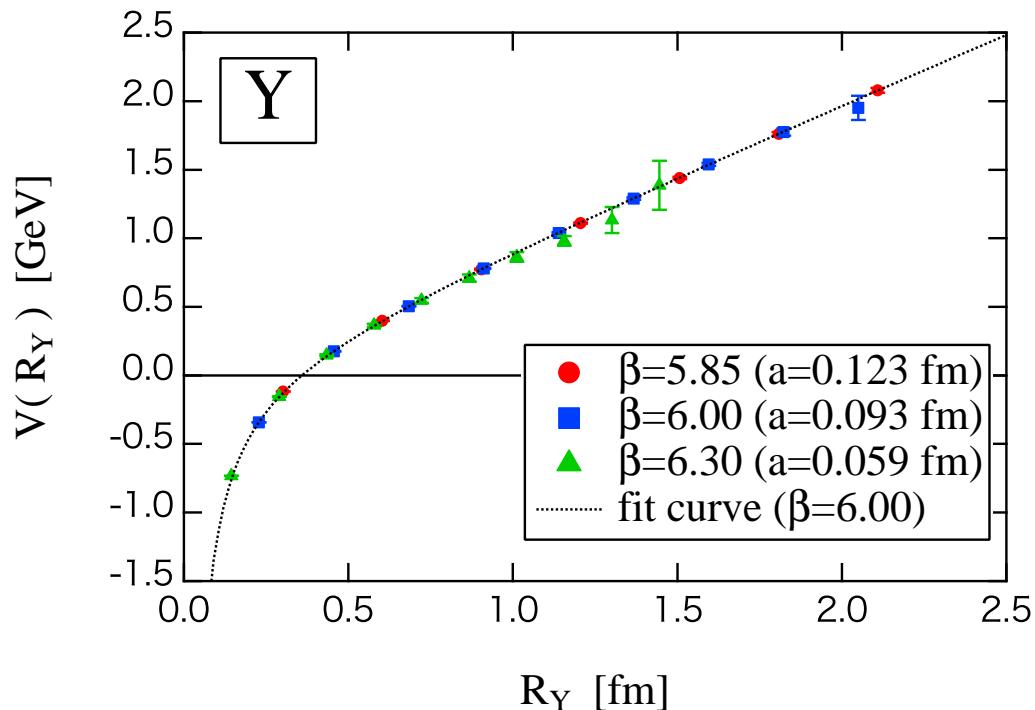
- ▶ Various 3-quark locations with  $Y$  and  $\Delta$  parametrization  
( $\beta=6.0$ ,  $24^4$  lattice,  $N_{\text{sub}}=6$ )



- ▶ “String tension” from  $Y$  parametrization is consistent with that of the  $Q\bar{Q}$  potential

# Three-quark potential & force (scaling test)

- Regular triangle with  $Y$  parametrization  
( $\beta = 5.85, 6.00, 6.30, 24^4$  lattice)



- The potential and force show good scaling behaviors

# Summary

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- ▶ We have investigated the static three-quark potential in SU(3) lattice gauge theory at zero temperature with the Polyakov loop correlation function (PLCF) by employing the multilevel algorithm
  - ⇒ PLCF allows to investigate the static potential of various 3-quark locations in detail with less systematic effects
  - ⇒ The potential with  $Y$ -parametrization falls into one curve whenever three quarks are separated with each other
  - ⇒ “String tension” from  $Y$ -parametrization is consistent with that of the  $Q\bar{Q}$  potential
  - ⇒ The potential and force show good scaling behaviors with  $Y$ -parametrization
- ▶ The present method is promising for computing relativistic corrections for the three-quark system